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# Self-Powered Piezoelectric and Thermoelectric Energy Simultaneous Extraction Interface Circuit Based on Double Stack Resonance

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**Abstract**—A hybrid self-powered synchronous electric charge extraction (HSP-SECE) interface circuit based on double stack resonance is presented in this paper. The proposed HSP-SECE interface circuit can simultaneously extract energy from piezoelectric transducer (PZT) and thermoelectric generator (TEG) when the peak open-circuit voltage of the PZT is detected by passive peak detector. The output power of the proposed interface circuit can reach three times of that of full-bridge rectifier circuit at the maximum power point, and the maximum efficiency of harvesting thermoelectric energy can reach 76% at 200 mV of the open-circuit voltage of the TEG. The simulation and experimental results show the superiority of the HSP-SECE circuit.

**Index Terms**—Double stack resonance, energy harvesting, piezoelectric transducer (PZT), self-powered, thermoelectric generator (TEG).

## I. INTRODUCTION

AS A BASIC component of the Internet of Things (IoT), wireless sensor network (WSN) nodes have been widely used in wearable devices, consumer electronics, and environmental monitoring. However, traditional battery-powered methods for WSN nodes have become an important obstacle to the development of WSNs. According to the new IoT technical standard published by the third Generation Partnership Project (3GPP) organization, the battery-powered life of WSN nodes under normal operation should be more than ten years [1]. Therefore, harvesting environmental energy is considered as one of the effective methods to overcome this challenge [2].

Piezoelectric energy harvesting (PEH) is one of the promising techniques that can scavenge energy from ambient vibration

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sources. With its electromechanical coupling characteristic, a piezoelectric element can generate electricity when strain is produced. Since the deformation in a vibrating structure is alternating, the generated electricity is also alternating [3]. Therefore, an interface circuit is needed between the piezoelectric transducer (PZT) and the load, which is of rectification, voltage regulation, and impedance matching [4]. The simplest interface circuit is a full-bridge rectifier (FBR) circuit, and many researchers have made their efforts in the optimization of the FBR and even implemented the circuit in complementary metal oxide semiconductor (CMOS) integration [5], [6]. However, its harvesting efficiency is low because the parasitic capacitance in the piezoelectric element causes a phase difference between the voltage and current and hence reactive power [7], [8].

Therefore, a variety of nonlinear energy extraction circuits have been presented to improve the energy harvesting efficiency. Lefeuvre *et al.* [9]–[11] proposed a parallel synchronized switch harvesting on inductor (P-SSHI) circuit, a synchronous electric charge extraction (SECE) circuit, and a series synchronized switch harvesting on inductor (S-SSHI) circuit. These circuits can effectively increase the harvesting efficiency, but require additional circuits to control the switches. Based on the S-SSHI circuit, Liang and Liao [3] proposed a self-powered synchronized switch harvesting on inductor (SP-SSHI) circuit. Compared to the FBR circuit, its maximum output power can increase by 200%. Based on the SECE circuit, Shi *et al.* [12] put forward a self-powered efficient SECE (SP-ESECE) circuit and its harvesting efficiency can reach up to 85.1%. Eltamaly and Addoweesh [13] proposed a self-powered synchronized switch harvesting on inductor (SP-SSHI) circuit developed from the P-SSHI circuit, which has 10% efficiency increase compared to the circuit in [3]. However, the reported circuits require large inductor to achieve good performance, which significantly increase the system volume. Although a flipping-capacitor rectifier was proposed to flip voltages using on-chip capacitors in [14], it is designed for high-frequency (>100 kHz) ultrasonic energy transfer applications but does not work for PZTs with large parasitic capacitance  $C_p$  [15].

Thermoelectric energy harvesting circuit mainly includes inductive boost converter and charge pump (CP) circuit. The main challenge of thermoelectric energy harvesting is that the open-circuit voltage of thermoelectric generator (TEG) is low, such as 236~265 mV at the temperature difference of 5 °C [16]. The

main limitations, indeed, come from the threshold voltage of diodes and pass transistors and the low transistor channel conductivity due to the low supply voltage [17]. Some researchers have proposed different methods, such as precharging output capacitors, adding mechanically assisted switches, and using external clocks or transformer startup circuits [18]. Das *et al.* [19] proposed a power-on-reset based self-starter with 2 nW quiescent power for thermoelectric energy harvester, which achieves a self-startup TEG voltage of 220 mV and a peak conversion efficiency of 76%. Yi *et al.* [20] used three-stage CP scheme employing a differential bootstrapped ring-VCO. The circuit has the efficiency of 38.8% at an input voltage of 150 mV and 45% at an input voltage of 200 mV. The CP circuit proposed by Kim *et al.* used dynamic body biasing and adaptive dead-time for efficiency improvement, which has an efficiency of 34% at an input voltage of 180 mV and 72.5% at an input voltage of 450 mV [21].

In order to harvest more energy, researchers have proposed some multisource harvesting circuits [22]–[25]. Lorenz *et al.* [22] presented a cooperative power harvesting scheme, collecting power from two or more independent ambient energy sources into a single nonlinear component. It is shown that the interaction of uncorrelated signals in a nonlinear device can greatly improve harvested power, but the conversion efficiency of the harvester is still low. Hsieh *et al.* [23] proposed a hybrid energy harvesting circuit in which both piezoelectric energy and solar energy are stored in a storage capacitor, and then an inductive boost circuit is used to power the load. However, the structure can only extract energy from energy source at higher open-circuit voltage. Yuk *et al.* [24] proposed a dual-source energy-harvesting interface circuit, which can harvest energy from both the PZT and the TEG. Yoon *et al.* [25] proposed a double pile-up resonance energy harvesting circuit that efficiently and simultaneously extracts energy from the PZT and the TEG. However, these circuits use an active control circuit module, which requires external power supply to start up the circuit.

In this paper, a hybrid self-powered SECE (HSP-SECE) interface circuit for harvesting dual-source energy is proposed. The proposed interface circuit can harvest both piezoelectric and thermoelectric energies simultaneously by double stack resonance.

This paper is organized as follows. In Section II, the circuit models of the PZT and the TEG, as well as the working principle of the FBR and the SECE circuits are introduced. In Section III, the architecture, working principle, and simulation analysis of the proposed interface circuit are described. In Section IV, the experimental results are provided. Section V concludes this paper.

## II. PRINCIPLES OF ENERGY HARVESTING

### A. Equivalent Model of the PZT

Fig. 1(a) shows the coupled model of the PZT [8]. In the mechanical domain,  $V_S$  represents the external excitation force,  $L_M$  represents the mechanical mass,  $C_K$  is the mechanical stiffness, and  $R_S$  takes into account the mechanical damping. In the

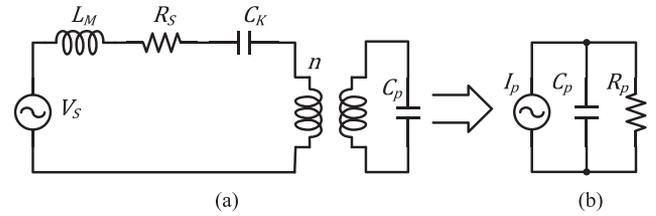


Fig. 1. PZT models. (a) Equivalent circuit model of the PZT. (b) Simplified circuit model under resonance conditions.

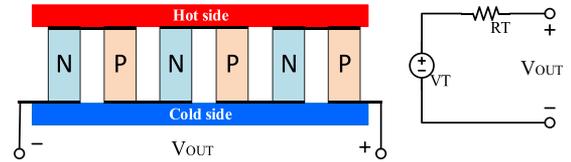


Fig. 2. Equivalent circuit of the TEG.

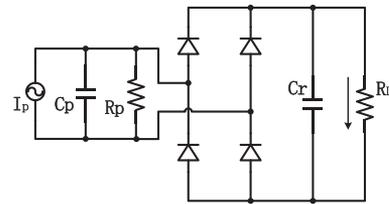


Fig. 3. Circuit of the FBR.

electrical domain,  $C_p$  represents the parasitic capacitance of the PZT. When the excitation is sinusoidal vibration, the vibration excitation of the PZT is modeled as a sinusoidal current source  $I_p$  in parallel with parasitic capacitance  $C_p$  and resistor  $R_p$ , as shown in Fig. 1(b) [26], [27].

### B. Equivalent Model of the TEG

Fig. 2 shows the structure of a typical TEG device. It consists of small legs of  $P$ - and  $N$ -type materials that are electrically connected in series. When a temperature difference is applied on two sides of the TEG, an open-circuit voltage between its two terminals is built because of the *Seebeck effect* [28]. In the literature, TEGs are mostly used under constant temperature gradient conditions. Moreover, the influence of the internal contact thermal resistance is usually neglected. Under these conditions, the TEG can be modeled as an equivalent constant voltage source in series with an equivalent internal resistance [29].

### C. Theoretical Analysis of the FBR

The FBR circuit is used to convert the alternating current (ac) of the PZT output into direct current (dc), as shown in Fig. 3.  $R_L$  is the equivalent resistance as the actual load, and the storage capacitance of  $C_r$  is much larger than the parasitic capacitance  $C_p$  of the PZT in order to reduce the output voltage ripple. When the open-circuit voltage  $V_{oc}$  of the PZT is greater than the load voltage  $V_{RL}$ , the rectifier bridge is turned ON and the energy stored in the PZT can be transferred to the load. Otherwise, the rectifier diodes are turned OFF.

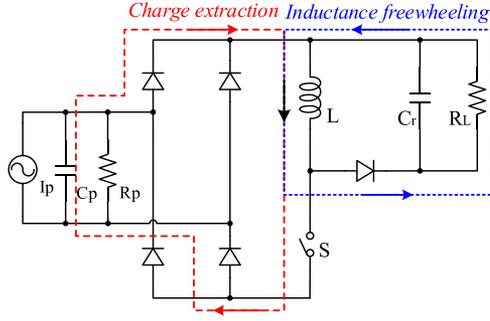


Fig. 4. Circuit of the SECE.

The output power of the FBR can be expressed as follows [30]:

$$P_{\text{FBR}} = f2 \int_0^{\frac{T_0}{2}} V_{RL} I_{\text{PZT}} dt = 2f\alpha V_{RL} \int_{\mu_1}^{\mu_M} du \quad (1)$$

where  $f$  is the vibration frequency of the PZT,  $T_0$  is the vibration period,  $\alpha$  is the force factor of the PZT,  $\mu_M$  is the peak displacement of the PZT, whereas  $\mu_1$  is the displacement of the PZT when the open-circuit voltage of the PZT is equal to the load voltage. The relationship between  $\mu_1$  and  $\mu_M$  can be expressed as

$$\mu_1 = \frac{2C_p V_{RL}}{\alpha} - \mu_M. \quad (2)$$

Hence, the output power of the FBR can also be expressed as

$$P_{\text{FBR}} = 4fV_{RL} [\alpha\mu_M - C_p (2V_D + V_{RL})]. \quad (3)$$

$V_D$  is the forward conduction voltage drop of diode.

The maximum output power and the corresponding optimal load resistance can be expressed as

$$P_{\text{FBR, max}} = fC_p \left( \frac{\alpha\mu_M}{C_p} - 2V_D \right)^2 \quad (4)$$

$$R_{\text{op}} = \frac{1}{4fC_p}. \quad (5)$$

If  $V_D$  is ignored, there is

$$P_{\text{FBR, max}} = f \frac{\alpha^2 \mu_M^2}{C_p} \quad (6)$$

#### D. Theoretical Analysis of the SECE

Compared to the FBR circuit, the SECE circuit adds an inductor  $L$ , a switch  $S$ , and a freewheeling diode, as shown in Fig. 4. The working principle of the SECE is that the charge accumulated in the parasitic capacitance  $C_p$  is transferred to the load side using the energy storage characteristics of the inductor. During one cycle, the switch is turned OFF at most of the time, the PZT is at an open state, and the energy in the equivalent current source  $I_p$  is transferred to the capacitor  $C_p$  [31].

The ideal current and voltage waveforms of  $C_p$  are shown in Fig. 5. When the switch is turned OFF, the equivalent current source  $I_p$  charges  $C_p$  so that the voltage across the PZT continuously rises. When the voltage is at its peak,  $I_p$  becomes

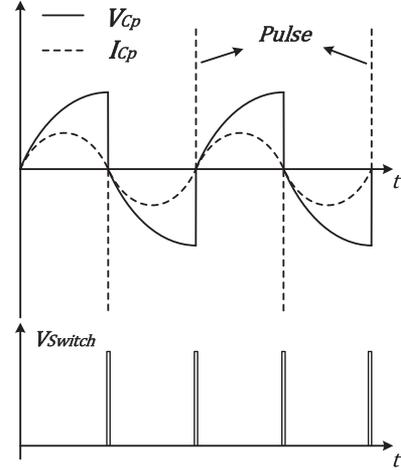


Fig. 5. Current and voltage of  $C_p$  and switching signal.

zero. At this time, the switch is turned ON so that the inductor  $L$  and capacitor  $C_p$  form a loop through the rectifying bridge, and hence the  $LC$  resonance occurs. After one-fourth resonance period, the switch is turned OFF immediately, and the accumulated energy on the  $C_p$  is completely transferred to the inductor. Meanwhile, the voltage quickly drops to zero. Then,  $L$  and  $C_p$  are disconnected, and the energy on the inductor can only be transferred to the load terminal through the freewheeling diode.

The  $LC_p$  loop resonance period can be expressed as

$$T_{LC_p} = 2\pi\sqrt{LC_p}. \quad (7)$$

Therefore, the turn-ON time of the switch can be expressed as

$$T_{S,\text{on}} = \frac{1}{4} T_{LC_p} = \frac{1}{2} \pi\sqrt{LC_p}. \quad (8)$$

Since  $C_p$  is small, the turn-ON time of the switch is very short. Hence, a narrow current pulse is generated on  $C_p$  when the inductor begins to extract the charge, as shown in Fig. 5.

The peak voltage of the PZT can be expressed as

$$V_{p,\text{max}} = \frac{2\alpha\mu_M}{C_p}. \quad (9)$$

Therefore, when the voltage reaches the peak point, the charge and energy accumulated on the capacitance  $C_p$  can be, respectively, expressed as

$$Q = C_p V_{p,\text{max}} \quad (10)$$

$$E_{C_p} = \frac{1}{2} C_p V_{p,\text{max}}^2 = \frac{2\alpha^2 \mu_M^2}{C_p}. \quad (11)$$

The principle of the SECE is to transfer the energy accumulated on  $C_p$  to the load twice in each cycle. Therefore, the output power of SECE can be expressed as

$$P_{\text{SECE}} = f2E_{C_p} = 4f \frac{\alpha^2 \mu_M^2}{C_p}. \quad (12)$$

According to the aforementioned analysis, the output power of the SECE circuit is constant, because the SECE periodically

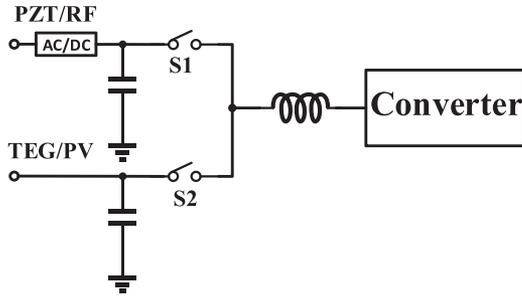


Fig. 6. Topological structure in [32] and [33].

extracts the energy accumulated on the parasitic capacitance through the  $LC$  resonance, which improves the energy harvesting efficiency and is not affected by the load. From (6) and (12), it can be found that the output power of the SECE is four times of the maximum output power of the FBR.

### III. PROPOSED INTERFACE CIRCUIT

For multisource energy harvesting, the main issue is that the extracted energy from high power source may flow back to the lower power one. The simple solution is to set a diode between the energy transducer and the converter. However, this has twofold defects: on one hand, part of the harvested energy may be consumed on it; on the other hand, when the open circuit voltage of the energy transducer is lower than the diode threshold voltage, no energy can be harvested.

There were some circuits proposed for multisource energy harvesting [32]–[33]. The principle of the topology structure is shown in Fig. 6. Although it can harvest energy from multisources, it has two defects. First, it ignores the optimization of the rectifier circuit for ac energy, such as from piezoelectric or RF transducer. In the traditional practice, the ac energy supplied to the load requires a two-stage conversion consisting of a rectification and a ac–dc conversion, in which each stage of conversion consumes part of the harvested energy. Second, in order to prevent energy backflow, the switches  $S_1$  and  $S_2$  cannot be turned ON simultaneously. Hence, maximum power extraction of two energy sources cannot be achieved. For achieving the high efficiency of multisource energy harvesting, these switches must be precisely controlled. Thus, a switching control module with external power supply is required, which cannot implement the energy harvester with self-starting.

To solve this problem, a topological structure for simultaneously harvesting ac and dc energies is proposed in this paper. As shown in Fig. 7,  $C_p$  is the storage capacitance of the ac current source, and  $C_t$  is the storage capacitor of the dc voltage source. When the voltage of  $C_p$  reaches a positive (negative) peak, the energy stored in the capacitor also reaches a maximum value. At this time, the switch  $S_p$  ( $S_n$ ) is turned ON, so that  $C_p$ ,  $C_t$ , and the inductor form an  $LC$  resonant circuit with double stack resonance. After one-fourth  $LC$  resonant period, the energy accumulated on the capacitors  $C_p$  and  $C_t$  is transferred to the inductor. Then, the switch  $S_p$  ( $S_n$ ) is immediately turned OFF while switch  $S_t$  is turned ON. The energy accumulated on the

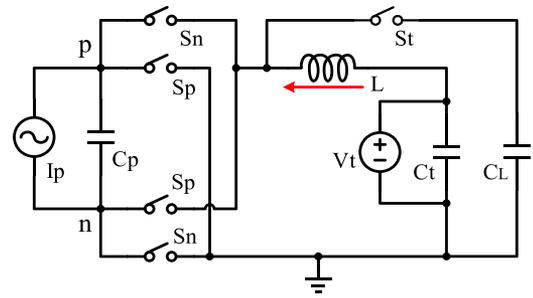


Fig. 7. Topological structure proposed in this paper.

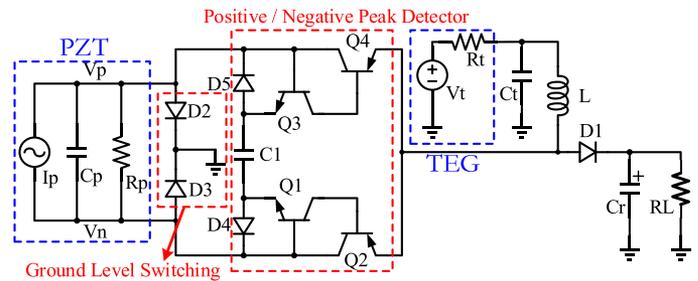


Fig. 8. Circuit of HSP-SECE.

inductor is transferred to the load. Thereby the simultaneous extraction of ac energy at the maximum power point and dc energy is achieved.

Based on the aforementioned structure, by introducing a ground-level switching circuit and a positive/negative peak detector circuit, a HSP-SECE circuit with double stack resonance is implemented, as shown in Fig. 8, which can simultaneously harvest the energy from the PZT and the TEG. It mainly includes the TEG, the PZT, positive/negative peak detector, ground-level switching circuit, capacitors  $C_r$ ,  $C_t$ , diode  $D_1$ , inductor  $L$ , and load resistor. Diode  $D_2$  and transistor  $Q_2$  are equivalent to switch  $S_p$  in Fig. 7, whereas diode  $D_3$  and transistor  $Q_4$  are equivalent to switch  $S_n$ , and diode  $D_1$  is equivalent to switch  $S_t$  in Fig. 7. Diodes  $D_4$  and  $D_5$ , transistors  $Q_1$  and  $Q_3$ , and capacitor  $C_1$  form a passive peak detector. Diodes  $D_2$  and  $D_3$  are turned ON in the positive/negative half cycle, respectively, so that the higher voltage end of the PZT is connected to the ground. Transistors  $Q_2/Q_4$  are turned ON at positive and negative peaks, respectively. The detail working principle is analyzed as follows.

#### A. Theoretical Analysis of Working Principle

Diodes  $D_2$  and  $D_3$  constitute a ground-level switching circuit connecting to the higher voltage end of the PZT to ground so that capacitors  $C_p$  and  $C_t$  and inductor  $L$  constitute a resonant circuit when the switch is turned ON. Positive/negative peak detector is to detect the peak of the open-circuit voltage. When the open-circuit voltage of the PZT reaches its peak, the switch  $Q_2$  (or  $Q_4$ ) is turned ON while capacitors  $C_p$  and  $C_t$  and inductor  $L$  form a resonant loop, so that the energy accumulated on capacitors  $C_p$  and  $C_t$  can be conveyed to inductor  $L$ , and then the energy is

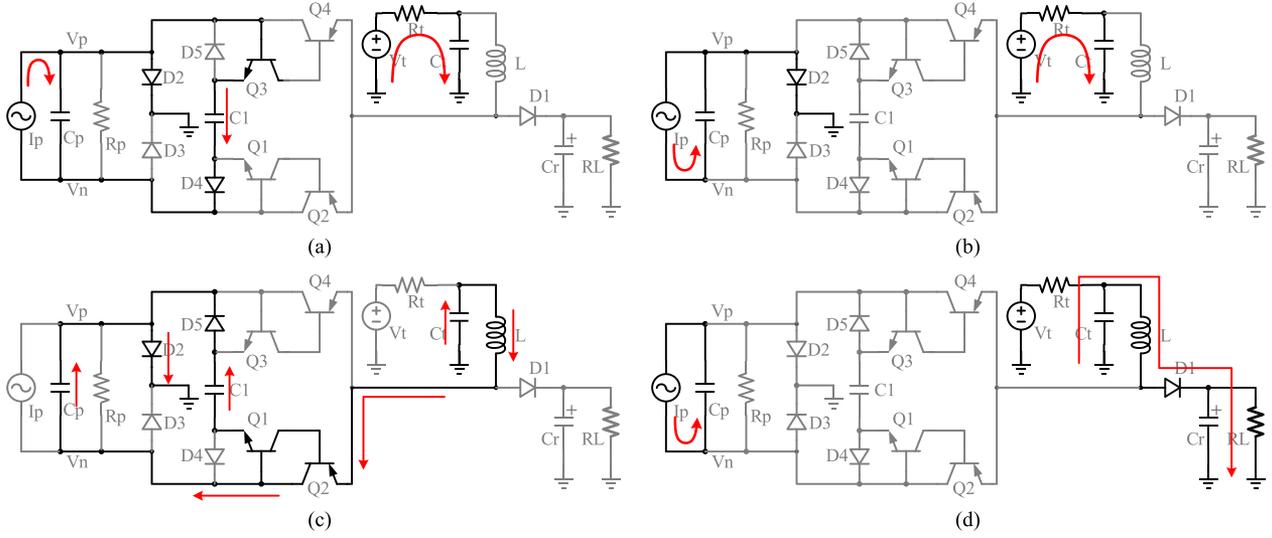


Fig. 9. Four phases within half of a vibration cycle. (a) Natural charging. (b) Current inversion. (c) Charge extraction. (d) Inductor freewheeling.

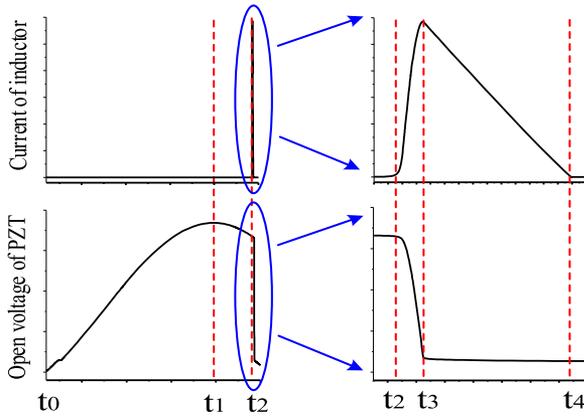


Fig. 10. Timing diagram of positive half cycle.

transferred to the load by freewheeling diode  $D_1$ . The specific working process is as follows.

In the positive half cycle, that is,  $V_p$  is higher than  $V_n$ , the circuit is mainly divided into four working phases, as shown in Fig. 9. The corresponding timing diagram is shown in Fig. 10, which includes the waveforms of the open-circuit voltage of the PZT and the current on the inductor at each working phase.

*Natural charging* ( $t_0 - t_1$ ): As shown in Fig. 9(a), the open-circuit voltage of the PZT gradually increases as the displacement increases. The parasitic capacitance  $C_p$  of the PZT and the capacitor  $C_1$  in the peak detection circuit are always in the charged state. At the same time, the TEG also charges capacitor  $C_t$ . When the open-circuit voltage of the PZT reaches its maximum value, the equivalent current source  $I_p$  just reaches zero, whereas the voltages on  $C_p$  and  $C_1$  both reach the maximum value. Then, the circuit starts to enter the second working phase.

When the open-circuit voltage of the PZT reaches an extreme value, the energy accumulated on  $C_p$  can be expressed as (11)

and the voltage across capacitor  $C_1$  also reaches the extreme

$$V_{C1,\max} = V_{p,\max} - V_{be} - V_D. \quad (13)$$

Hence, the energy accumulated on  $C_1$  can be expressed as

$$E_{C1} = \frac{1}{2} C_1 V_{C1,\max}^2 = \frac{1}{2} c_1 (V_{p,\max} - V_{be} - V_D)^2. \quad (14)$$

*Current inversion* ( $t_1 - t_2$ ): As shown in Fig. 9(b), when the PZT begins to move in a reverse direction, the equivalent current source  $I_p$  starts to charge  $C_p$  inversely, so that the open-circuit voltage of the PZT starts to decrease. Due to the threshold voltage of diode  $D_5$ , the charge on  $C_1$  is unable to be released and the voltage across  $C_1$  remains unchanged. When the voltage across  $C_1$  is higher than the open-circuit voltage of the PZT by the sum of the threshold voltage of transistor  $Q_1$  and the turn-ON voltage drop of diode  $D_5$ , transistor  $Q_2$  is turned ON while the voltage across capacitor  $C_t$  reaches its maximum value. Then, the circuit begins to enter the third working phase.

The energy accumulated on  $C_t$  can be expressed as

$$E_{Ct} = \frac{1}{2} C_t V_{Ct,\max}^2. \quad (15)$$

The open-circuit voltage of the PZT at this time is

$$V_p' = V_{C1,\max} - V_{be} - V_D = V_{p,\max} - 2V_{be} - 2V_D. \quad (16)$$

The corresponding energy accumulated on  $C_p$  is

$$E'_{Cp} = \frac{1}{2} C_p V_p'^2 = \frac{1}{2} C_p (V_{p,\max} - 2V_{be} - 2V_D)^2. \quad (17)$$

Hence, there is a phase lag, which can be expressed as  $\theta$

$$\theta = \cos^{-1} \left( \frac{V_{p,\max} - 2V_{be} - 2V_D}{V_{p,\max}} \right). \quad (18)$$

The energy loss caused by the phase lag can be expressed as

$$\Delta E_\theta = E_{Cp} - E'_{Cp} = \frac{1}{2} C_p V_{p,\max}^2 \sin^2 \theta. \quad (19)$$

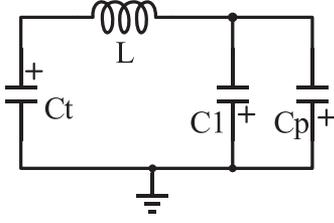


Fig. 11. Equivalent circuit of charge extraction.

It can be seen that the decrease of the extracted energy is closely related to  $\theta$ . The efficiency of energy extraction can be expressed as

$$\eta = \frac{E_{C_p} - \Delta E_\theta}{E_{C_p}} = \cos^2 \theta. \quad (20)$$

The smaller the phase lag  $\theta$ , the higher the efficiency of energy extraction.

**Charge extraction ( $t_2 - t_3$ ):** When the transistor  $Q_2$  is turned ON,  $C_p$ ,  $C_1$ ,  $C_t$ , and  $L$  constitute an  $LC$  resonant loop, as shown in Fig. 9(c). If the turn-ON voltage drops of the transistor and diode are ignored, the equivalent circuit is shown in Fig. 11. It is noted that the high voltage terminals of  $C_1$  and  $C_p$  are connected to ground, which is implemented through the ground-level switching circuit. Since the capacitance value of  $C_1$  is much smaller than  $C_p$ ,  $C_1$  can be ignored. The equivalent circuit can be regarded as a single  $LC$  resonant circuit composed of  $C_p$ ,  $C_t$ , and the inductor  $L$ , which is called double stack resonance. After one-fourth  $LC$  resonance period, the charges stacked on  $C_t$  and  $C_p$  are transferred to the inductor  $L$ . Therefore, the charge extracted by the inductor is increased compared to the pure SECE circuit.

The energy transferred to the inductor can be calculated as

$$E'_L = E'_{C_p} + E_{C_t} + E_{C_1}. \quad (21)$$

According to (21), compared to published self-powered PEH interface circuits, the HSP-SECE circuit can not only harvest  $E'_{C_p}$  but also simultaneously harvest thermoelectric energy  $E_{C_t}$  based on double stack resonance.

However, in the charge extraction process, both the diode and the transistor consume energy as long as there are currents flowing through them. As shown in Fig. 9(c), there are currents flowing through four devices,  $Q_2$ ,  $D_2$ ,  $Q_1$ , and  $D_5$ . However, the current through  $Q_2$  and  $D_2$  is much larger than that through  $Q_1$  and  $D_5$ . Hence, the energy consumed by the former is much larger than the later. Assuming that the energy consumption by the former is considered and the turn-ON voltage drops of  $Q_2$  and  $D_2$  remain unchanged, the energy consumed in the charge extraction process can be expressed as

$$E_{CE,loss} = \int (V_D + V_{CE}) I(t) dt. \quad (22)$$

Here  $V_D$  and  $V_{CE}$  represent the voltage drops across the diode and transistor, respectively.

**Inductor freewheeling ( $t_3 - t_4$ ):** After the energy on the capacitor  $C_1$  is fully discharged,  $Q_2$  is turned OFF. Through the freewheeling diode  $D_1$ , the inductor  $L$  and the load form a loop

TABLE I  
DATA OF THE COMPONENTS

Component	Value
$I_p$	300 $\mu$ A/20Hz
$C_p$	200nF
$R_p$	2M $\Omega$
$C_1$	2.8nF
$V_t$	200mV
$R_t$	10 $\Omega$
$C_t$	100 $\mu$ F
$L$	0.5mH
$C_r$	200 $\mu$ F
$RL$	50k $\Omega$

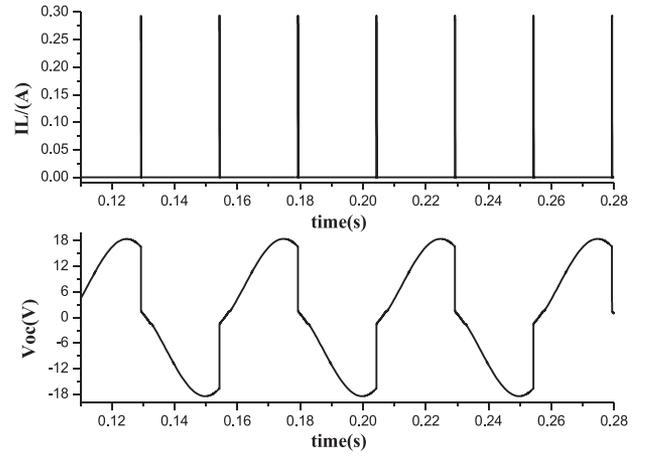


Fig. 12. Current of  $L$  and open-circuit voltage of PZT.

to transfer the energy accumulated on the inductor to the load, and the equivalent current source starts to charge the  $C_p$  inversely, as shown in Fig. 9(d). Then, the energy extraction of the positive half cycle is completed.

During the process of the inductor freewheeling, the components also consume energy. If the internal resistance of the inductor is neglected, the main energy loss causes from the freewheeling diode  $D_1$  is

$$E_{IF,loss} = \int V_D I(t) dt. \quad (23)$$

From another perspective, the conversion efficiency of the inductor freewheeling process can be simplified as

$$\eta_{IF} = \frac{V_{out}}{V_{out} + V_D}. \quad (24)$$

From (24), it can be seen that the conversion efficiency increases with output voltage increase.

During the negative half cycle, the HSP-SECE circuit works in a similar way.

## B. Simulation Analysis

The simulation of the proposed HSP-SECE circuit is performed. The parameters of the components used are shown in Table I. Fig. 12 shows the PZT open-circuit voltage and current through the inductor  $L$ . It can be seen that both the inductor current and the open-circuit voltage of the PZT are periodically

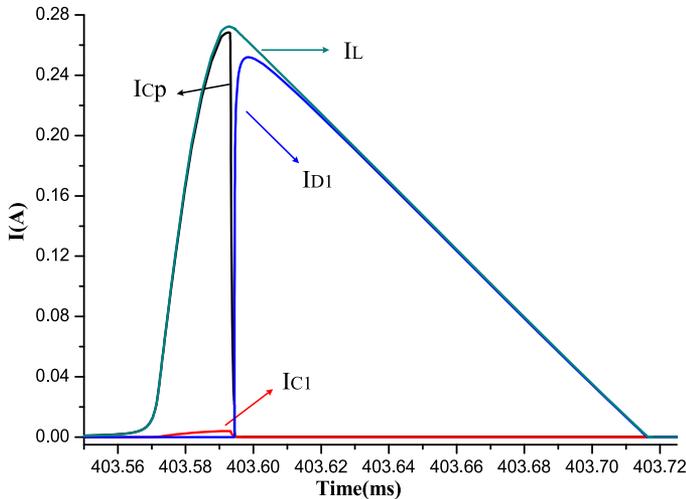


Fig. 13. Current waveforms of  $L$ ,  $C_p$ ,  $C_1$ , and  $D_1$ .

varied. When the inductor generates a current pulse, the open-circuit voltage of the PZT drops rapidly. It can be seen that there is a phase lag between the current pulse and the peak voltage of the PZT, which confirms (18).

In order to understand the energy extraction process of the HSP-SECE circuit, the relevant current waveform is zoomed in to observe the detail clearly at the instant of switching conduction. Fig. 13 shows the currents on inductor  $L$ , capacitors  $C_p$  and  $C_1$ , and freewheeling diode  $D_1$ , respectively. According to the previous theoretical analysis, before the current peak, it is the energy extraction process, whereas after the current peak, it is the inductor free-wheeling process. When the inductor current rises, the currents through capacitors  $C_p$ ,  $C_1$ , and  $C_t$  have the same direction and the sum of the currents of the capacitors  $C_p$  and  $C_1$  equals the inductor current, which confirms the analysis for Fig. 11 and in other words, the energy accumulated on capacitor  $C_1$  can be recycled. As the inductor current reaches its peak, the currents of capacitors  $C_p$  and  $C_1$  drop to zero. At this point, the energy extraction process ends. When the energy on the capacitor  $C_1$  is completely released, the peak detector turns the switch ( $Q_2$  or  $Q_4$ ) OFF, and the circuit immediately enters the inductor free-wheeling process. Then, the current of diode  $D_1$  rises quickly to the inductor current, which means that the energy extraction is completed. However, near the peak of the diode current, there is a difference between the inductor current and the diode current. This is that the circuit uses a passive peak detector to control the switching transistor  $Q_2$  (or  $Q_4$ ), so that the transistor cannot be quickly disconnected, which leads to partial current leakage.

### C. Comparative Analysis

In order to compare the load dependence of several different circuit structures and the advantages of each type of the circuits, the HSP-SECE circuit proposed in this paper and the published circuits are simulated. The simulation is based on the equivalent circuit model shown in Fig. 1 and the value of the load resistance is adjusted while the parameters  $I_p$ ,  $C_p$ ,  $R_p$ , and load

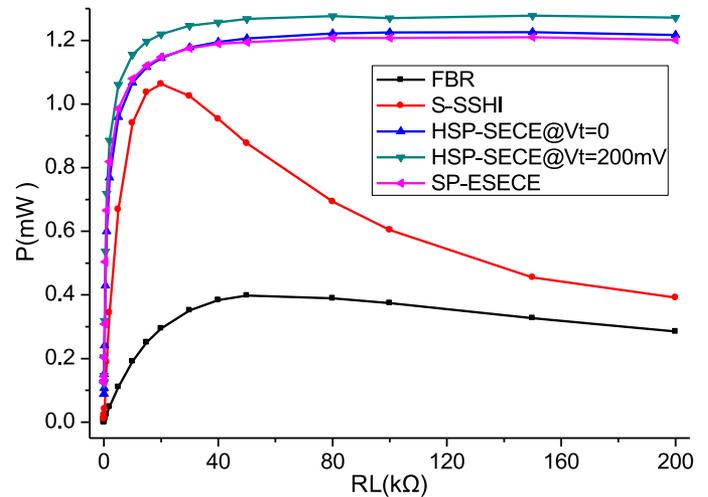


Fig. 14. Output power comparison of several circuits under the same excitation and different loads.

capacitance  $C_r$  keep constant. The constant parameters are the same as Table I. Since a purely resistive load is used, according to *Ohm's law*, the output power varies with the load, which also means that the output power changes with the output voltage. Fig. 14 shows the output power of the FBR, S-SSHI circuits in [3], SP-ESECE circuit in [12], and the proposed HSP-SECE circuit at  $V_t = 0$  mV and  $V_t = 200$  mV, respectively. From the simulation results, the maximum output power of the HSP-SECE, SP-ESECE, and S-SSHI circuits are 3.05 ( $V_t = 0$  mV), 3, and 2.64 times of that of the FBR circuit, respectively.

From Fig. 14, it can be seen that the output power of the S-SSHI and SP-ESECE circuits is higher than that of the FBR circuit in the entire load range. It also can be seen that when the load resistance is lower than 20 k $\Omega$ , the output power is small, and increases as the load value increases, which is in accordance with the description of (24). However, when the load resistance exceeds 20 k $\Omega$ , the output power of the S-SSHI circuit begins to decrease significantly as the load increases, whereas the output power of the SP-ESECE and the HSP-SECE circuits remains almost unchanged. This also proves that the harvesting efficiency of the SECE circuit is hardly affected by the load.

When the load is higher than 50 k $\Omega$ , the output power of the HSP-SECE circuit is higher than that of the SP-ESECE circuit even when the TEG open-circuit voltage is 0 mV, and it is significantly higher than that of the S-SSHI circuit. For the HSP-SECE circuit, when the TEG open-circuit voltage is 200 mV, the output power is significantly higher than that when the TEG open-circuit voltage is zero. It shows that the HSP-SECE circuit can effectively extract the thermoelectric energy even if the open-circuit voltage of the TEG is as low as 200 mV.

In order to test the advantages of the HSP-SECE circuit for harvesting thermoelectric energy, the circuit under the same piezoelectric excitation with different output voltages is simulated. The simulation results are shown in Fig. 15. Three sets of data are measured, corresponding to TEG's open circuit

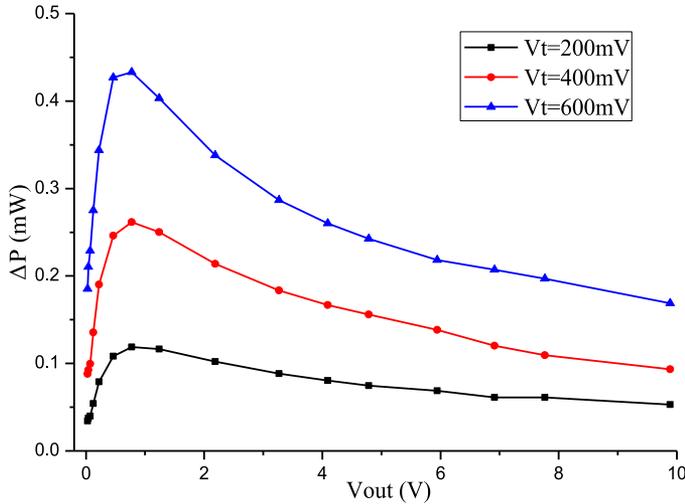


Fig. 15.  $\Delta P$  under different output voltage.

voltages of 200, 400, and 600 mV, respectively. In the figure, the abscissa represents the effective value of the output voltage of the circuit, whereas the ordinate  $\Delta P$  stands for the difference between the output power under the specific  $V_t$  and the output power when  $V_t = 0$  mV. From Fig. 15, it can be seen that when the output voltage is very low,  $\Delta P$  is small, which means that the advantage of the hybrid energy harvester is not obvious. This is mainly that most of the energy is consumed by the freewheeling diode when the output voltage is low, which is consistent with (24). Therefore, the output power will gradually increase as the output voltage increases. When the output voltage is about 1 V,  $\Delta P$  reaches its maximum value. However, as the output voltage increases,  $\Delta P$  begins to decrease slowly. This is that during the inductor freewheeling process, as the output voltage increases, the anode voltage of the freewheeling diode gradually increases, and the positive voltage of the capacitor  $C_t$  increases. Therefore, the current charged by the TEG to capacitor  $C_t$  is reduced so that the output power of the TEG slightly decreases.

#### IV. EXPERIMENTAL WORK

In order to verify the aforementioned theoretical and simulation analyses, experimental platform is set up, as shown in Fig. 16. The experimental system is composed of a function signal generator, an oscilloscope, a PZT, a dc source, a power amplifier, a vibrating table, a TEG device, and the HSP-SECE energy harvesting circuit. The sinusoidal signal generated by the signal generator is used to control the vibrating table after being enhanced by the power amplifier. One end of the PZT is tightly fixed with a metal mass, and the other end is fixed on the vibrating table. When the vibrating table works, due to the inertia of the metal mass, a relative displacement between two ends of the PZT occurs so that the PZT deforms accordingly and generates an ac signal. Before connecting the circuit, the vibration frequency and amplitude need to be adjusted by the signal generator and power amplifier. Only when the vibration frequency is adjusted to the resonance state, the open circuit

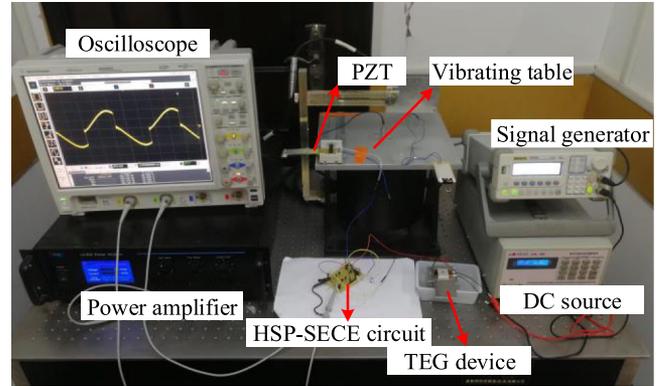


Fig. 16. Experimental setup.

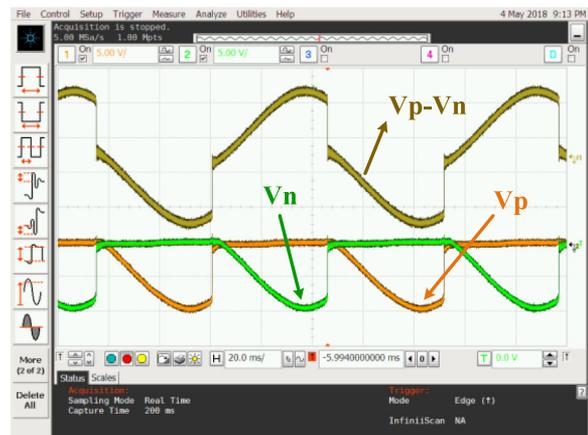


Fig. 17. Experimental waveforms.

voltage of the PZT reaches the maximum. The TEG device consists of a TEG, a container filled with cold water, a positive temperature coefficient (PTC) heating sheet and a heat sink. The heated surface of the TEG is attached to the PTC heating sheet, whereas the cold surface is placed on the heat sink, which is placed in the cold water container. The PTC heating sheet is powered by a dc source. Adjusting the output voltage of the dc source can control the temperature of the heating sheet, which can change the temperature difference across the TEG. When the output voltages of the TEG and PZT are stable, they are connected to the circuit for testing.

Fig. 17 shows the voltage waveforms measured by the oscilloscope, which corresponds to the voltage waveforms at both ends of the PZT and the open-circuit voltage waveforms, respectively. When the PZT's open-circuit voltage reaches its peak, it rapidly drops to zero after a delay, the voltage starts to reverse, and the PZT is at an open-circuit state. Therefore, the PZT's open-circuit voltage begins to increase with the increase of the PZT's displacement, which is in line with the simulation results.

Fig. 18 shows the measured output power of the HSP-SECE circuit when the excitation frequency of the PZT is maintained at 22 Hz and the load resistance remains at 150 k $\Omega$ . The temperature difference between the two sides of the TEG and the PZT vibration amplitude is gradually increased. It can be seen

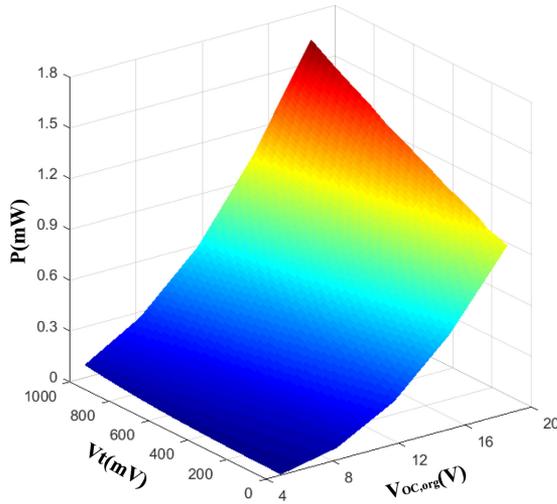


Fig. 18. Output power at different open-circuit voltages of TEG and PZT (with  $C_t = 100 \mu\text{F}$ ).

that the output power of the circuit is positively correlated with the amplitude of the PZT and the temperature difference of the TEG. When the amplitude of the PZT remains unchanged, the output power tends to increase linearly with the increase of the open circuit voltage of the TEG. When the temperature difference between the two ends of the TEG remains unchanged, the output power increases exponentially with the increase of the original open-circuit voltage of the PZT. This indicates that the HSP-SECE circuit can harvest energy from the PZT and the TEG at the same time. Furthermore, it can be seen that as the open-circuit voltage of the PZT increases, the curvature of the curve increases significantly with the thermoelectric voltage increase, which means that as the amplitude of the vibration increases, the harvesting efficiency of the thermoelectric energy can be enhanced.

In order to verify whether TEG has an effect on the efficiency of PEH, we tested the output power of the HSP-SECE circuit when the open-circuit voltage of the TEG was 0 V and the TEG was shorted (equivalent to an SECE circuit for PEH) when  $V_{OC,org} = 13 \text{ V}$  remained unchanged, as shown in Fig. 19. As illustration, when  $C_t$  is small, the TEG significantly impair the efficiency of PEH. While  $C_t \geq 60 \mu\text{F}$ , the output power of the circuit is almost the same as that of the SECE circuit. Therefore, in practical applications, it is necessary to select a larger value of  $C_t$ , so that the efficiency of PEH is not affected by the TEG.

Table II summarizes comparison results among the HSP-SECE circuit and those state-of-the-art circuits, which are for PEH in [3] and [12] based on S-SSHI and SECE, respectively, thermoelectric energy harvesting in [20] and [21] based on charge-pump and both piezoelectric energy and thermoelectric energy harvesting in [24] and [25].

Compared with those single energy harvesting methods, the HSP-SECE circuit is of comparative advantages. In terms of PEH, the harvesting efficiency of the HSP-SECE can be up to three times of that of the FBR circuit with only a 0.5-mH inductor though the maximum efficiency of the S-SSHI in [3] can

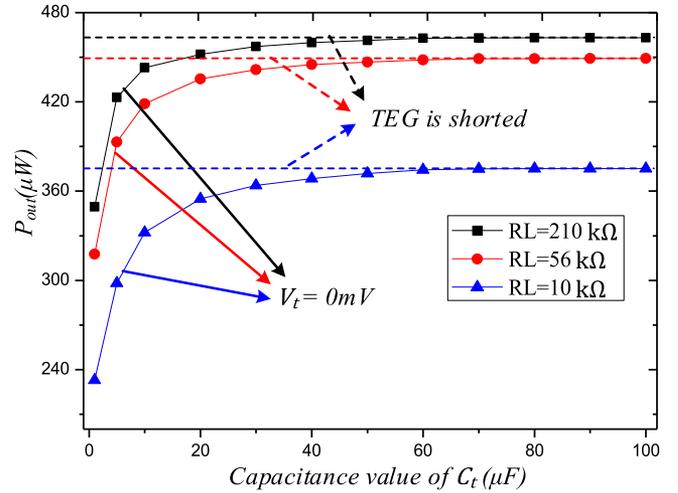


Fig. 19. When  $V_t = 0 \text{ mV}$ , the output power changes with  $C_t$ .

also reach three times of the FBR circuit with a 47-mH inductor and under optimal load. For thermoelectric harvesting, on one hand, the harvesting efficiency of the HSP-SECE is as high as 76% at input voltage 200 mV, whereas the efficiencies from the state-of-the-art circuits in [20] and [21] are 45% at input voltage 200 mV and 72.5% at input voltage 450 mV, respectively, and their harvesting efficiencies are rapidly decreased as the input voltages decrease; on the other hand, the minimum starting voltage of the HSP-SECE circuit can be as low as 50 mV, whereas that in [20] and [21] is 150 mV.

In terms of dual energy-source harvesting, from Table II, it can be seen that the HSP-SECE circuit not only can simultaneously harvest piezoelectric and thermoelectric energies but also is self-powered. The circuits in [24] and [25] from the same research group can also simultaneously harvest piezoelectric and thermoelectric energies. It can be seen that the maximum efficiencies of thermoelectric harvesting for both the HSP-SECE circuit and the circuit in [25] have similar value. Although the PEH efficiency of the HSP-SECE is only three times of that of the FBR, whereas the circuits in [24] and [25] can be as high as 4.22 and 14.52 times, respectively, the latter requires an external power supply and hence the power loss of the circuit is hard to be calculated. Once the battery is exhausted, the harvester will be paralyzed and unable to restart. In addition, since the circuits in [24] and [25] are implemented in CMOS process, the piezoelectric energy at low open-circuit voltage can be harvested, so that the maximum output power can only be as low as 87 and 16  $\mu\text{W}$ , respectively. The output power of the HSP-SECE circuit is less affected by the load size and maintains a high harvesting efficiency over a wide load range. However, the circuit in [25] can maintain a high harvesting efficiency only in the optimal load region. If the load is too large or too small, the efficiency will be significantly reduced. The circuit in [25] contains a boost circuit, which can separately harvest thermoelectric energy. The circuit proposed in this paper can harvest piezoelectric energy separately, but it is not possible to separately harvest thermoelectric energy.

TABLE II  
PERFORMANCE COMPARISON

Publication	TIE 2012[3]	JIMSS 2016[12]	TCS II 2018 [20]	JSSC 2015[21]	ISSCC 2014[24]	JSSC 2018[25]	This Work
Input Sources	PZT	PZT	TEG	TEG	PZT+TEG	PZT+TEG	<b>PZT+TEG</b>
PZT	N/A	N/A	N/A	N/A	N/A	PPA-1001	<b>PPA-1011</b>
Features	S-SSHI	SECE	Charge-Pump	Charge-Pump	Pile-Up Resonance	Double Pile-Up Resonance	<b>Double Stack Resonance</b>
Self-Power	YES	YES	YES	YES	NO	NO	<b>YES</b>
Inductor	47mH	1mH	N/A	N/A	N/A	0.68mH	<b>0.5mH</b>
Boost to FBR <sub>(max)</sub>	3× @1.5mW	3× @2mW	N/A	N/A	4.22× @87μW	14.52× @16μW	3× <b>@1mW</b>
Peak efficiency of TEG	N/A	N/A	38.8%@150mV 45%@200mV	34%@180mV 72.5%@450mV	N/A	75%(booster)	<b>76%@200mV</b>
Startup voltage of TEG	N/A	N/A	150mV	150mV	500mV	100mV	<b>50mV</b>
Frequency/Hz	30	50	N/A	N/A	100	140	<b>22</b>

## V. CONCLUSION

In this paper, the method to simultaneously harvest the energy from PZT and TEG was studied. An HSP-SECE circuit based on double stack resonance was proposed. Theoretical analysis, simulation, and experimental results showed that an HSP-SECE circuit can effectively harvest both piezoelectric and thermoelectric energies simultaneously, and maintain a high efficiency for ultralow voltage thermoelectric energy harvesting. However, it should be pointed out that for the proposed circuit, when the open-circuit voltage of the TEG was zero, the internal resistance of the TEG might consume a small part of the power harvested from the PZT.

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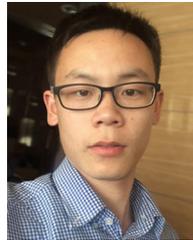
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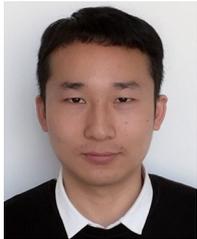
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